A silver metal spiral binding is visible on the left side of the notebook cover, consisting of a series of loops that hold the pages together.

The RF Photoinjector Gun: Microwave point-of-view

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RF gun as microwave device

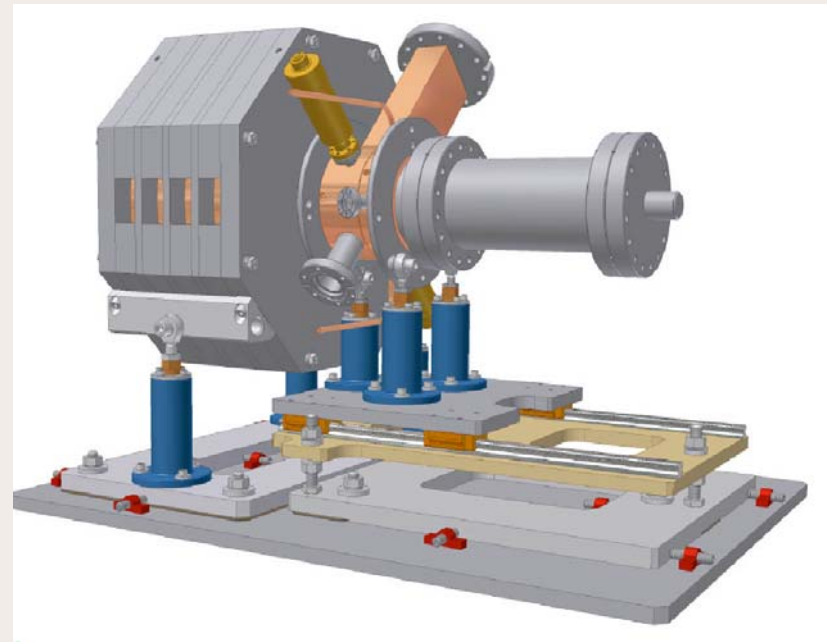
- Photoinjector guns are high gradient standing wave devices
 - Beam dynamics complex
 - Space-charge and RF are large effects
- Need to understand (a lot!):
 - Cavity resonances
 - Coupled cavity systems
 - Power dissipation
 - External coupling
 - Time-domain response
 - Measurement and tuning procedures
- Reference: Chapter 7 *Fundamentals of Beam Physics*

The photoinjector gun



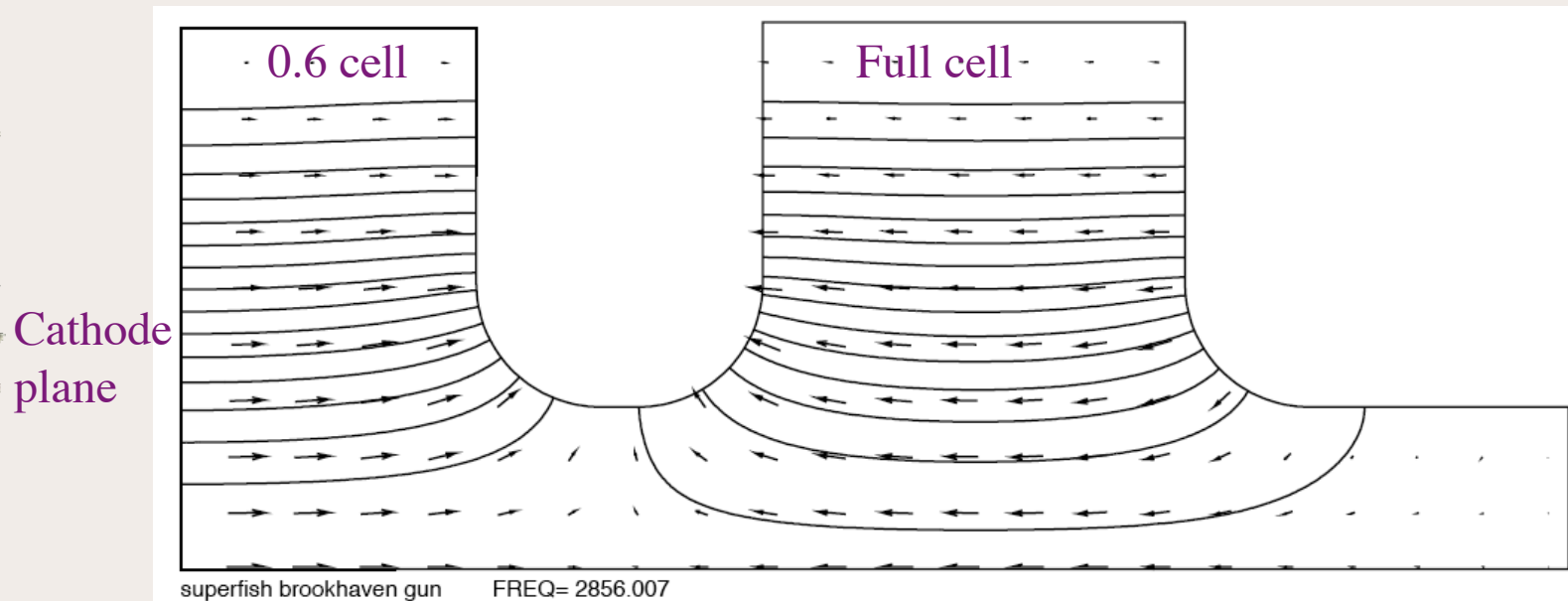
Neptune RF gun
with cathode plate removed

SPARC gun and solenoid



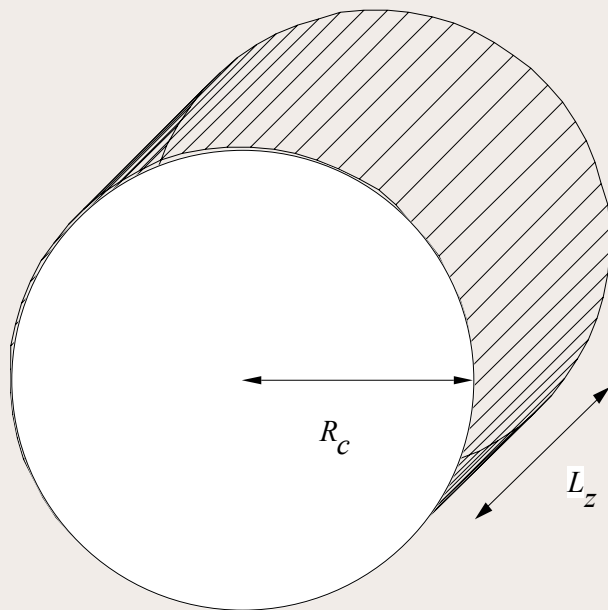
The 1.6 cell rf gun geometry

- Approximate with cylindrical symmetry



- π -mode, full ($\lambda/2$) cell with 0.6 cathode cell
- How to calculate resonant frequencies?

Cavity resonances: the pill-box model

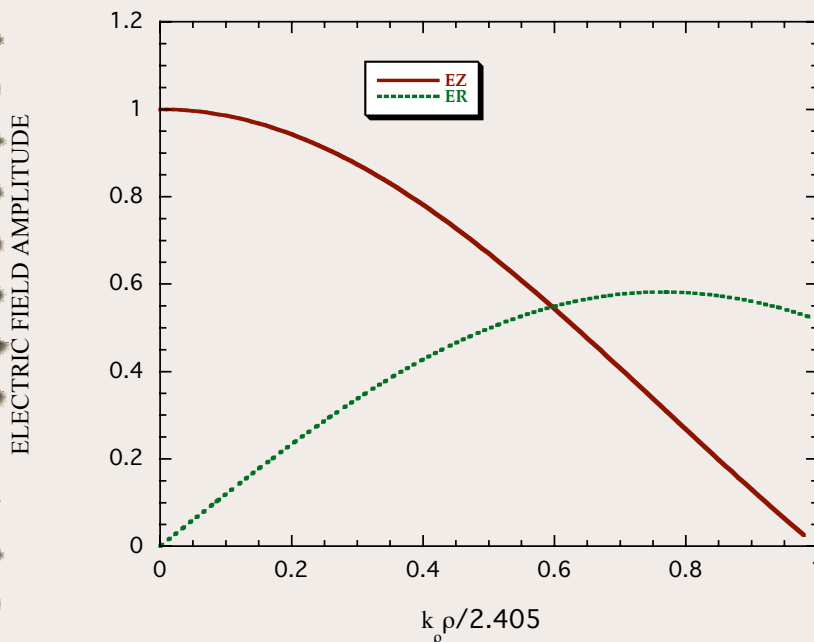


- Pill-box model approximates cylindrical cavities
- Resonances from Helmholtz equation analysis

$$\left[\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial}{\partial \rho} \right) - k_{z,n}^2 + \frac{\omega^2}{c^2} \right] \tilde{R} = 0$$

- Look for TM (axial field) modes
- Conducting BCs

Pill-box fields



- No longitudinal dependence in fundamental

$$k_{z,0} = 0 \quad \omega_{0,1} \cong \frac{2.405c}{R_c}$$

- Electromagnetic fields

$$E_z(\rho) = E_0 J_0(k_{\rho,0}\rho)$$

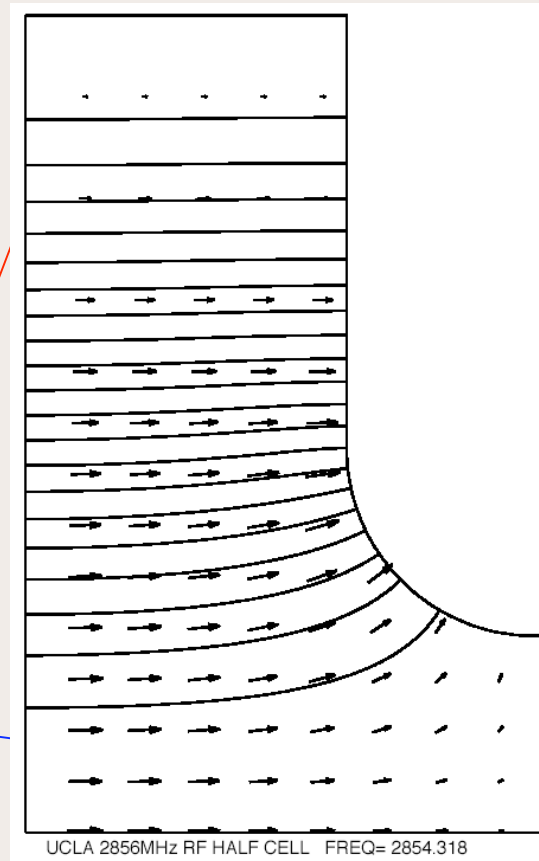
$$H_\phi(\rho) = \frac{\omega \epsilon_0}{k_{\rho,0}} E_0 J_1(k_{\rho,0}\rho) = c \epsilon_0 E_0 J_1(k_{\rho,0}\rho)$$

- Stored energy

$$\begin{aligned} U_{EM} &= \frac{1}{4} \epsilon_0 L_z E_0^2 \int_0^{R_c} [J_0^2(k_{\rho,0}\rho) + J_1^2(k_{\rho,0}\rho)] \rho d\rho \\ &= \frac{1}{2} \epsilon_0 L_z E_0^2 R_c^2 J_1^2(k_{\rho,0} R_c) \end{aligned}$$

A circuit-model view

- Lumped circuit elements may be assigned: L , C , and R .
- Resonant frequency
$$\omega \cong \frac{1}{\sqrt{LC}}$$
- Tuning by changing inductance, capacitance
- Power dissipation by surface current (H)



Contours of constant flux in 0.6 cell of gun

Cavity shape and fields

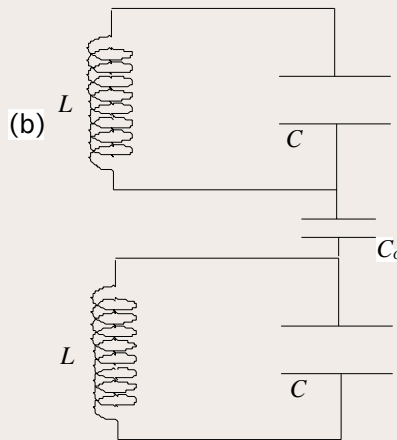
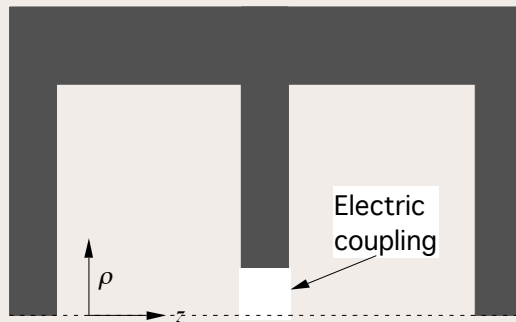
- Irises needed for beam passage and *RF coupling*
- Fields *near axis* (in iris region) may be better represented by spatial harmonics

$$E_z(\rho, z, t) = E_0 \operatorname{Im} \sum_{n=-\infty}^{\infty} a_n \exp[i(k_{n,z}z - \omega t)] I_0[k_{\rho,n}\rho] \quad k_{\rho,n} = \sqrt{k_{n,z}^2 - (\omega/c)^2}$$

- Higher (no speed of light) harmonics have nonlinear (modified Bessel function) dependence on ρ .
 - Energy spread
 - Nonlinear transverse RF forces
- Avoid re-entrant nose-cones, etc. Circle-arc sections are close to optimum.

Cavity coupling: simple model

(a)



- Circuit model allows simple derivation of mode frequencies
- Ignore resistance in walls (does not affect frequency much)
- Electric coupling is *capacitive*

$$\frac{d^2 I_1}{dt^2} + \omega_0^2(1 - \kappa_c)I_1 = -\kappa_c \omega_0^2 I_2$$

$$\frac{d^2 I_2}{dt^2} + \omega_0^2(1 - \kappa_c)I_2 = -\kappa_c \omega_0^2 I_1$$

- Off-axis slots can give magnetic coupling (effectively reverses sign in κ)
- Solve eigenvalue/eigenmode equations

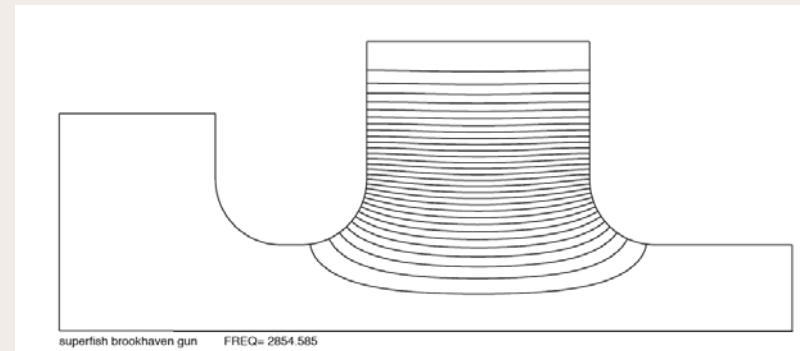
$$\begin{bmatrix} \omega_0^2(1 - \kappa_c) - \omega^2 & \kappa_c \omega_0^2 \\ \kappa_c \omega_0^2 & \omega_0^2(1 - \kappa_c) - \omega^2 \end{bmatrix} \mathbf{i} = 0$$

Coupled cavity modes

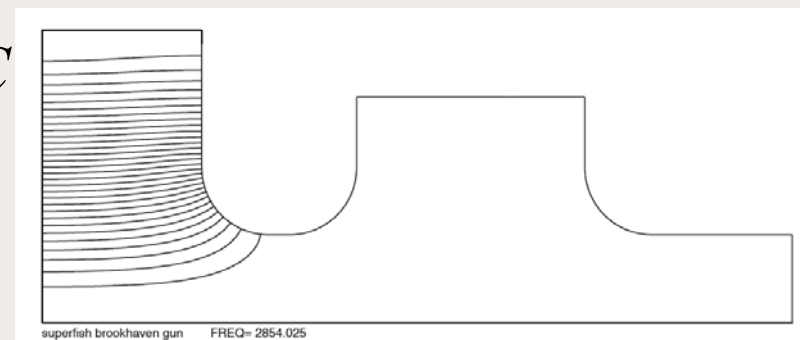
- Secular equation $\omega^4 - 2\omega_0^2(1 - \kappa_c)\omega^2 + \omega_0^4(1 - 2\kappa_c) = 0$
- Two eigenvalues $\omega = \omega_0$ (0 - mode) and $\omega = \omega_0\sqrt{1 + 2\kappa_c}$ (π - mode)
- Corresponding to eigenvectors $\mathbf{i} = \frac{1}{\sqrt{2}} \begin{pmatrix} \pm 1 \\ 1 \end{pmatrix}$
- Lower frequency when cells are in phase (0-mode)
- The higher frequency mode is the π -mode, where the excitation is 180 degrees out of phase in the two cavity cells
 - Higher frequency due to less near-axis fields
 - Lower capacitance

Finding frequencies of cells

- Real geometry, use SUPERFISH simulation
- In simulation detune the opposing cell
- 0.6 cell has frequency of 2854.01 MHz
- Full cell is slightly higher: 2854.60 MHz due to lower C
- These are the geometries which produce “field” balance



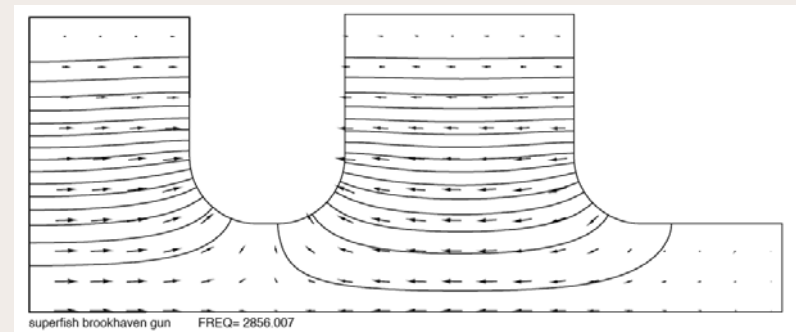
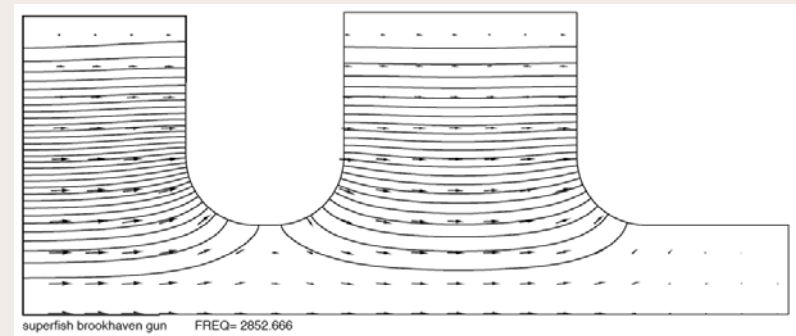
Full cell



0.6 cell

The coupled mode frequencies

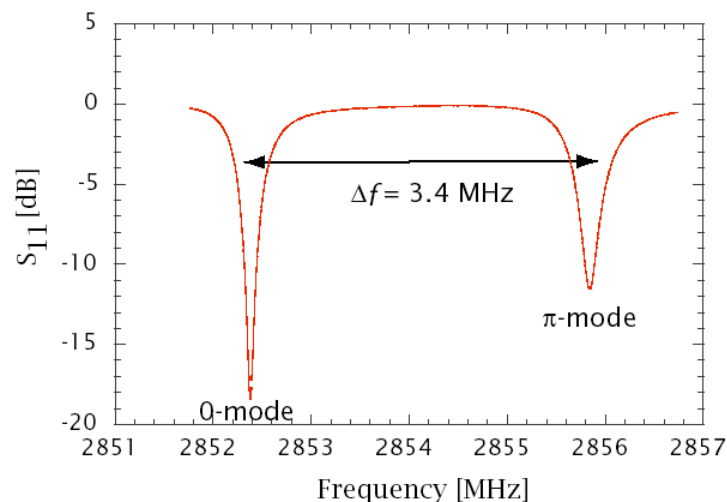
- The π -mode frequency is 2856.0 MHz as expected
- This mode needs to be balanced (equal fields)
- Zero mode frequency is 2852.66 MHz. Note that field arrows do not reverse
- The frequency separation is ~ 3.34 MHz according to calculation
 - This is sensitive to the iris geometry as built



Measurement and tuning of frequencies

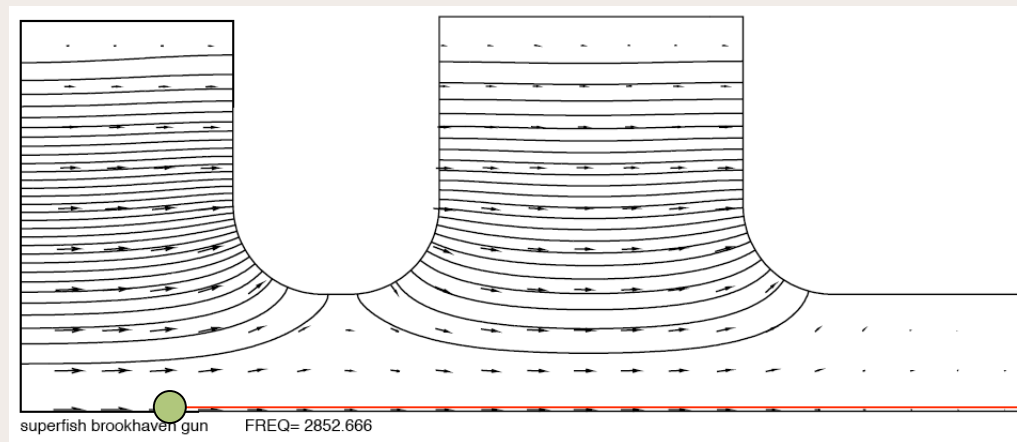
- Frequency response can be measured on a network analyzer
 - *Calibrated* phase and amplitude response, calibrated frequency
- Two measurements: S_{11} (reflection) S_{21} (transmission)
- Resonance frequencies of individual cells and coupled modes
 - Full cell. Remove cathode to detune 0.6 cell.
 - 0.6 cell. S_{11} through (detuning) probe in full cell.
- Tuning via Slater's theorem/circuit model guide

$$\frac{\delta\omega_0}{\omega_0} = \frac{\delta V_c}{U_{EM}} \left[\frac{1}{2} \epsilon_0 \vec{E}^2 - \frac{1}{2} \mu_0 \vec{H}^2 \right]$$

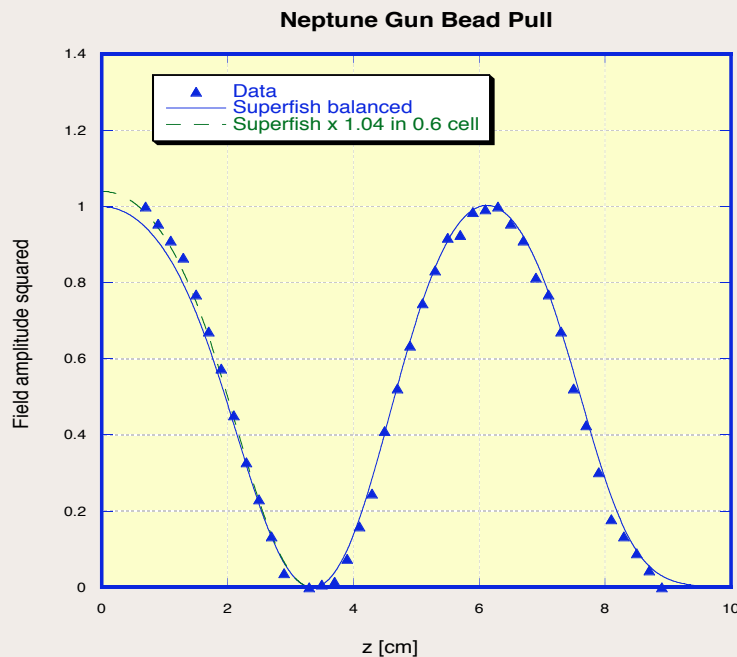


Measurement of fields

- Use so-called “bead-pull” technique
 - “Bead drop in gun, beam tube pointed upwards...
- Metallic or dielectric bead (on optical fiber)
- Metallic bead on-axis gives negative frequency shift (electric field energy displaced)
 - No magnetic effects on-axis for accelerating mode $|E_z| \propto \sqrt{-\delta\omega}$
- More complex if one has magnetic fields (deflector)

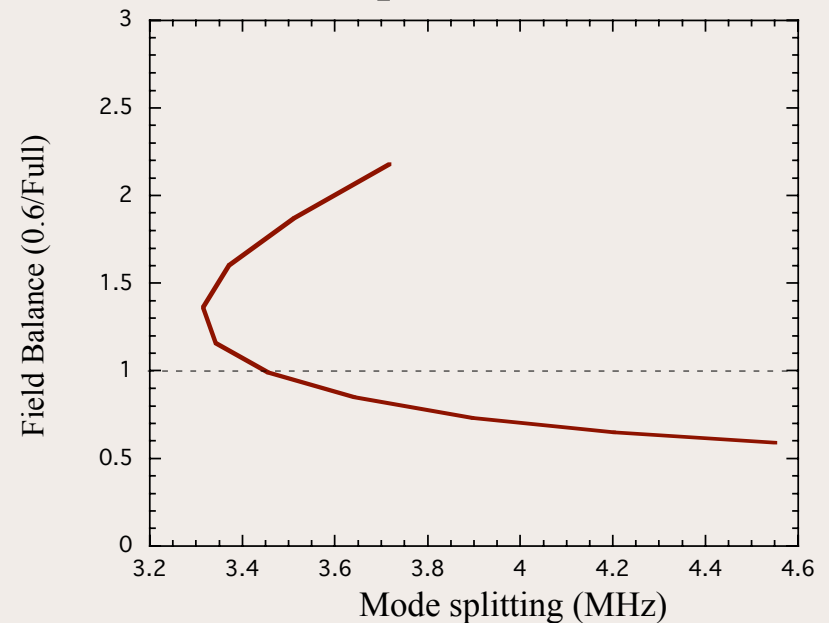


Field balance and mode separation

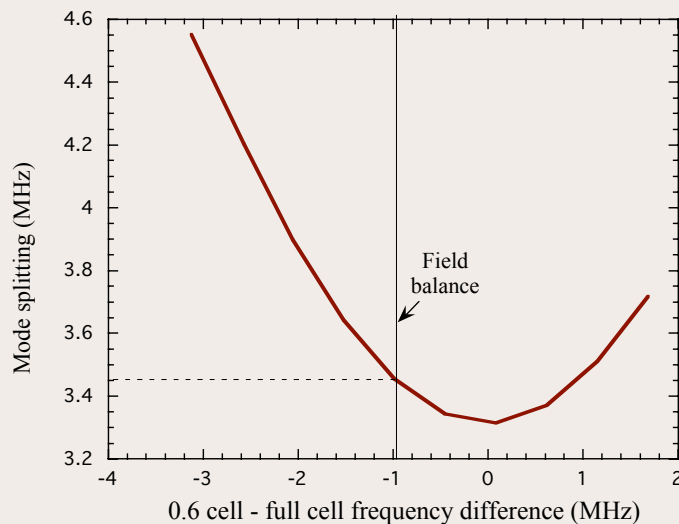
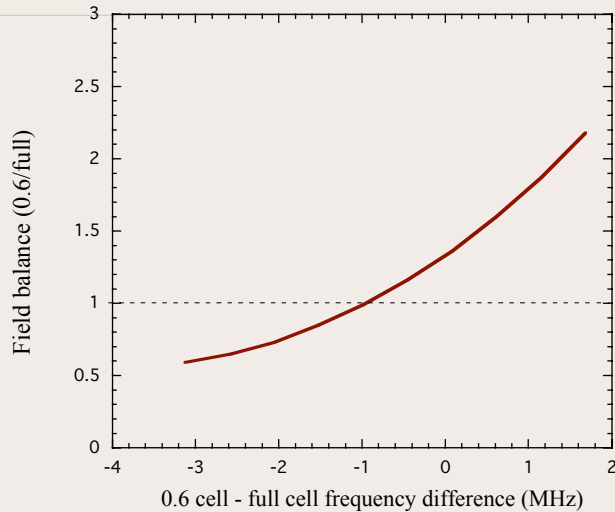


New Neptune gun field balance from bead pull

- Bead-pull technique is invasive
- Calibrate mode separation
 - Slightly different than coupled cavity model...
- Field balance possible in fast cathode swap



Other parameterizations

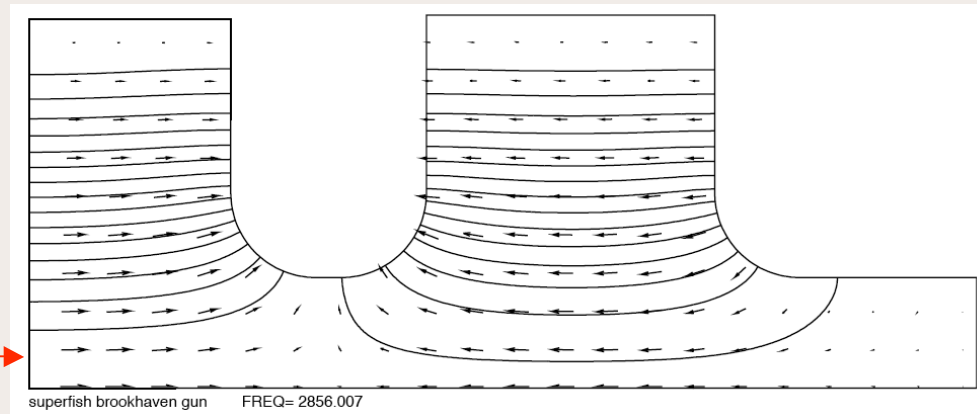


- Fundamental problem: the field balance is a double valued function of the mode splitting
- Full cell frequency is tuned by insertable (magnetic) tuners (in=higher)
- If you find minimum splitting, make full cell higher in frequency

Final tuning

- Do not want the full-cell tuners re-entrant (breakdown)
- Cathode deformation using “tuning nut”
- Temperature tuning=44 kHz/degree
- Almost perfect balance between atmospheric index and 20° C above room temperature operation

Capacitive
tuning thru
deformation →



Temperature
tuning

Power dissipation in walls

- Wave equation in conductors, harmonic solution

$$\left[\vec{\nabla}^2 - \mu_0 \sigma_c \frac{\partial}{\partial t} - \mu_0 \epsilon \frac{\partial^2}{\partial t^2} \right] \begin{Bmatrix} \vec{E} \\ \vec{B} \end{Bmatrix} = 0 \quad -k^2 + i\omega\mu_0\sigma_c + \mu_0\epsilon\omega^2 = 0$$

- Complex wavenumber into wall normal, gives skin-depth

$$k = \sqrt{\frac{\omega\mu_0\sigma_c}{2}}(1+i) \quad \delta_s = [\text{Im}k]^{-1} = \sqrt{\frac{2}{\omega\mu_0\sigma_c}}$$

- Power is lost in a narrow layer (skin-depth) of the wall by surface current excitation $K_s = |\vec{H}_{\parallel}| = \mu_0 |\vec{B}_{\parallel}|$

$$\frac{dP}{dA} = -\frac{K_s^2}{4\delta_s\sigma_c} = -\frac{K_s^2}{4} \sqrt{\frac{\omega\mu_0}{2\sigma_c}} = -\frac{K_s^2}{2} R_s, \quad R_s \equiv \frac{1}{2} \sqrt{\frac{\omega\mu_0}{2\sigma_c}} \quad \text{Surface resistivity}$$

Power dissipation and Q

- Total power dissipated in walls (pill-box example)

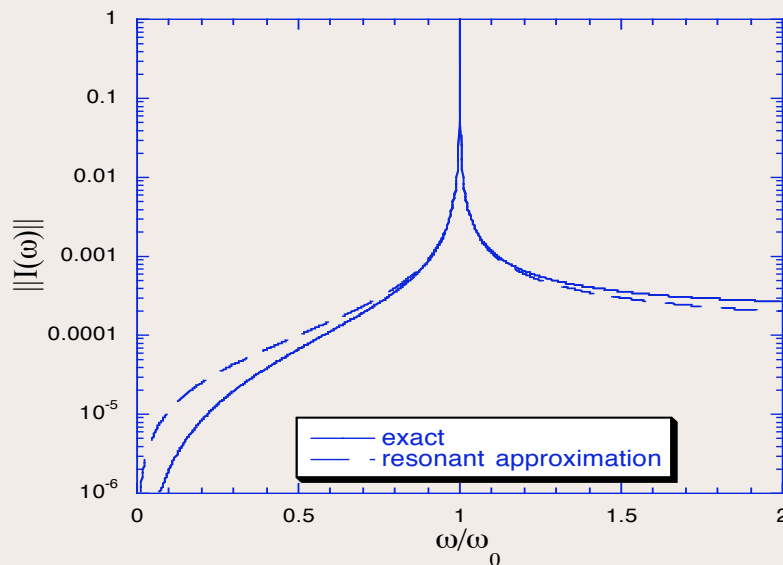
$$\begin{aligned}\langle P \rangle &= \pi R_s (c \epsilon_0 E_0)^2 \left[R_c L_z J_1^2(k_{\rho,0} R_c) + 2 \int_0^{R_c} J_1^2(k_{\rho,0} \rho) \rho d\rho \right] \\ &= \pi R_s (c \epsilon_0 E_0)^2 R_c J_1^2(k_{\rho,0} R_c) [L_z + R_c]\end{aligned}$$

- Internal quality factor (pill box)

$$Q \equiv \frac{\omega U_{EM}}{\langle P \rangle} = \frac{Z_0}{2R_s} \frac{2.405 L_z}{(R_c + L_z)} \quad Z_0 = 377 \Omega$$

- For ~ 2856 MHz (S-band), $Q \sim 12,000$
- Other useful interpretations of Q
 - Exponential response in time domain
 - Frequency response half-width

Frequency domain picture



$$Q = \frac{\omega}{\Delta\omega_{1/2}} = \frac{L}{R}$$

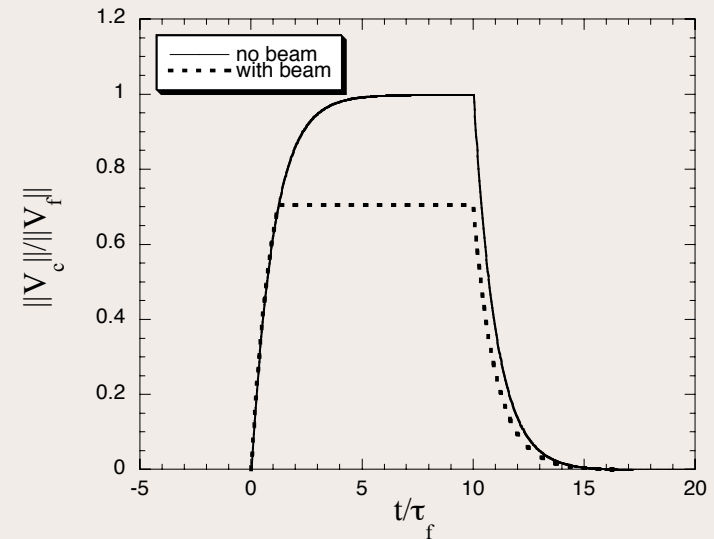
- Resonant width proportional to inverse of power loss
- Allows definition of lumped resistance
- Can measure with network analyzer
- Be careful about external coupling

Cavity filling, emptying, and external coupling

- Exponential response of cavity voltage

$$V_c(t) \cong \frac{2\beta_c}{1 + \beta_c} V_F \sin([\omega_0(t - t_0)]) \left[1 - \exp\left(-\frac{\omega_0}{2Q_L}(t - t_0)\right) \right]$$

- Controlled by *loaded* Q_L
- Energy extraction by
 - Radiation into waveguide
 - Beam acceleration (high average current systems)



VSWR and β

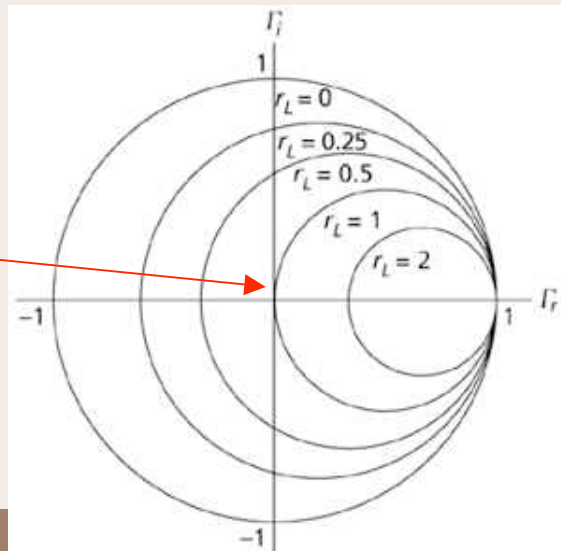
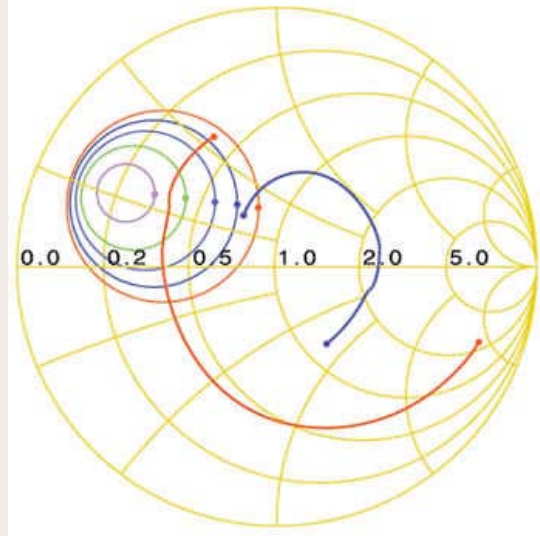
- Measure reverse and forward voltage
 - Time domain
 - NWA Smith chart

- Calculate VSWR

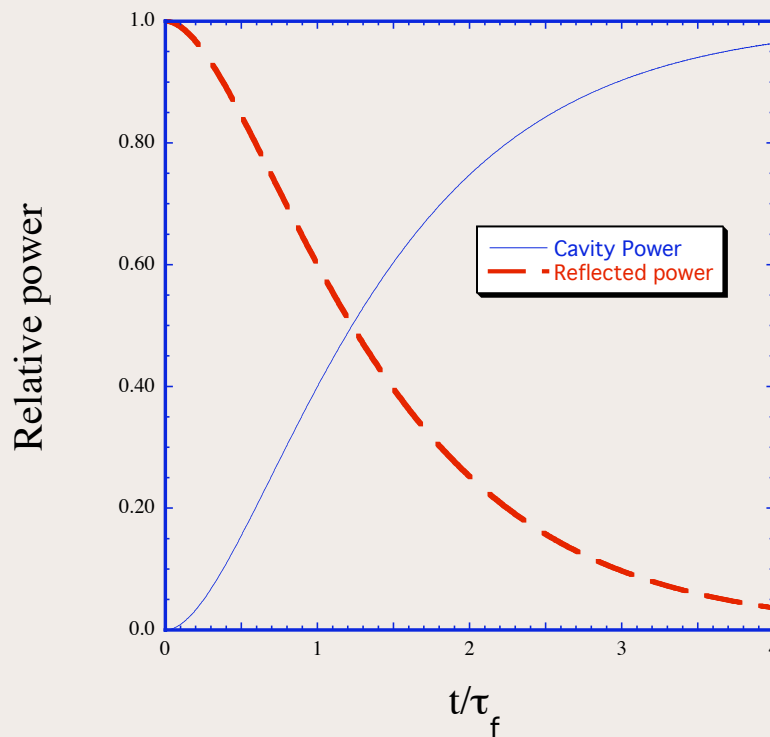
$$VSWR \equiv \frac{\|V_F\| + \|V_R\|}{\|V_F\| - \|V_R\|} = \begin{cases} \beta_c, & \beta_c > 1 \\ \beta_c^{-1}, & \beta_c < 1. \end{cases}$$

- We want critical coupling $\beta_c = 1$
- This is $r_c = 1$ on lower right

- Loaded Q $Q_L = \frac{Q}{1 + \beta_c} = \frac{Q}{2}$



Temporal response of the cavity



- Standing wave cavity fills exponentially
$$E \propto 1 - \exp\left(-\frac{\omega t}{2Q}\right) \propto 1 - \exp\left(-\frac{t}{\tau_f}\right)$$
- Gradual *matching* of reflected and radiate power (E^2) from input coupler
 - Reflected wave from input coupler=re-radiated wave
- In steady-state, all power goes into cavity (critical coupling) so reflected power is eventually cancelled

Voltage, power and acceleration

- The square of the cavity voltage is proportional to the shunt impedance

$$P_c = \frac{\omega U_c}{Q_0} = \frac{L \|I_c\|^2}{\omega L / R} = \|I_c\|^2 R = \frac{\|V_c\|^2}{R} = \frac{V_c^2}{2R} = \frac{V_c^2}{Z_s}$$

- The accelerating field is The square of the cavity voltage is proportional to the shunt impedance per unit length

$$Z'_s = \frac{dP}{dz} / \langle E_0^2 \rangle$$

- The 1.6 cell gun S-band cavity has

$$Z'_s \cong 40 \text{ M}\Omega/\text{m}$$

The 1.6 cell RF gun

- The design power for the Neptune gun is $P \sim 6$ MW
- The “accelerating” length of the gun is $L_g = 0.0845$ m
- The average accelerating field is

$$E_0 = \sqrt{P \frac{Z'_s}{L_g}} = 53 \text{ MV/m}$$

- The peak on-axis accelerating field is \sim twice the average, or over 100 MV/m